

Feasibility of measuring the electromagnetic moments of Λ_c at the LHC

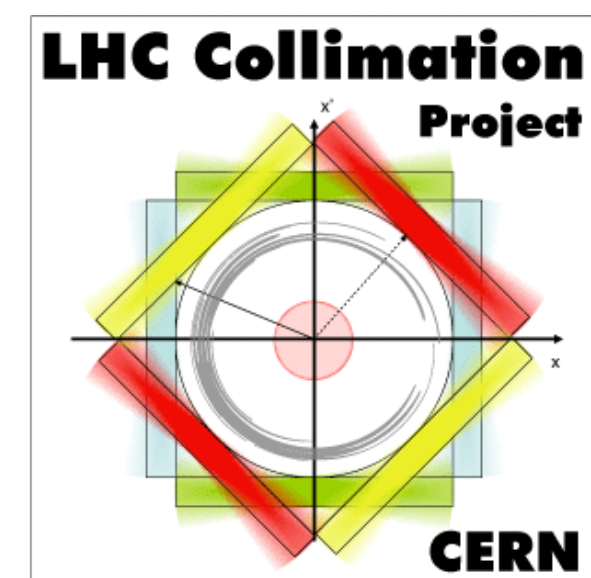
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Introduction

- Electromagnetic moments of baryons
- Spin precession in a bent crystal

Optimal crystal orientation for EDM measurement

- Spin precession in a bent crystal
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- Quantitative analysis [1]

Performance assessment of layouts in IR3 and IR8 of LHC

- Double crystal layouts at LHC [3,4]
- Precision of measurement [1]
- Possible improvements [1,3]

[1] [A.S. Fomin et al. Eur. Phys. J. C \(2020\) 80:358](#)

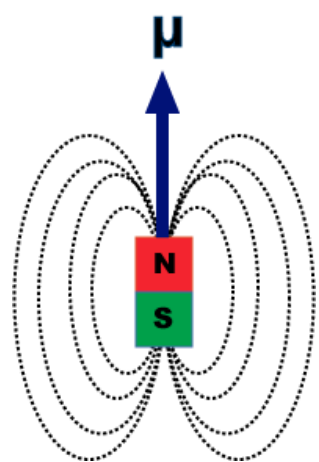
[2] [A.S. Fomin, JHEP 08 \(2017\) 120](#)

[3] [D. Mirarchi et al. Eur. Phys. J. C 80 \(2020\) 10, 929](#)

[4] [CERN Yellow Reports: Monographs, 4/2020](#)

Electromagnetic moments of baryons

Magnetic Dipole Moment:

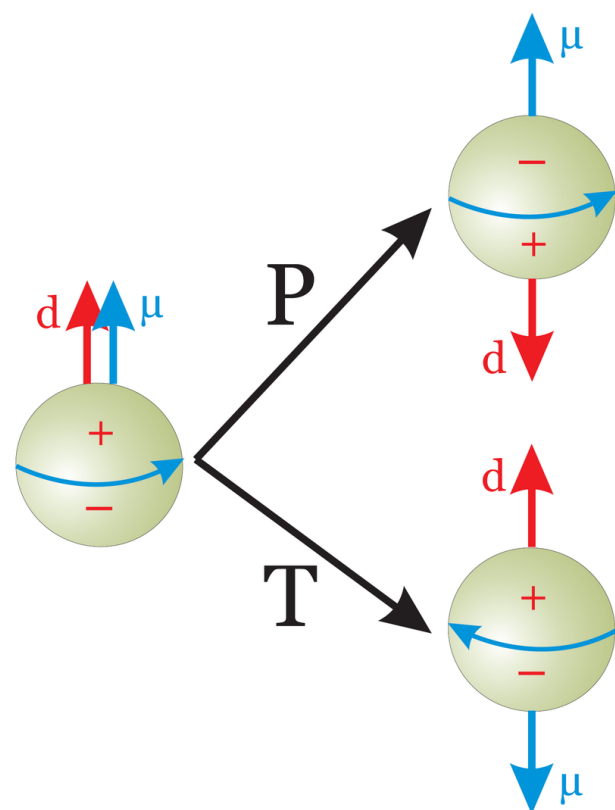


$$\vec{\mu} = \frac{g}{2} \frac{e}{m} \vec{S}, \quad \vec{S} = \frac{\hbar}{2} \vec{\sigma}$$

- $|g| = 2 \rightarrow$ a point-like Dirac particle
- $|g| \approx 2 \rightarrow$ a radiative corrections
- $|g| \neq 2 \rightarrow$ a composite structure or NP

Particle	CT	g -factor	Comments	
p	∞	+ 5.585 694 702 (17)	exp.	
n	$\sim \infty$	- 3.826 085 45 (90)	exp.	
Σ^+	2.4 cm	+ 6.233 (25) + 6.1 (1.2) _{stat} (1.0) _{syst}	exp.	world-average value using Bent Crystals (at Fermilab 1990)
Λ_c^+	60 μ m	+ 1.90 (15) not measured	theor.	assuming $g_c \approx 2$
			exp.	Feasibility studies at LHC

Electric Dipole Moment:



$$\vec{\delta} = \frac{f}{2} \frac{e}{m} \vec{S}, \quad \vec{S} = \frac{\hbar}{2} \vec{\sigma}$$

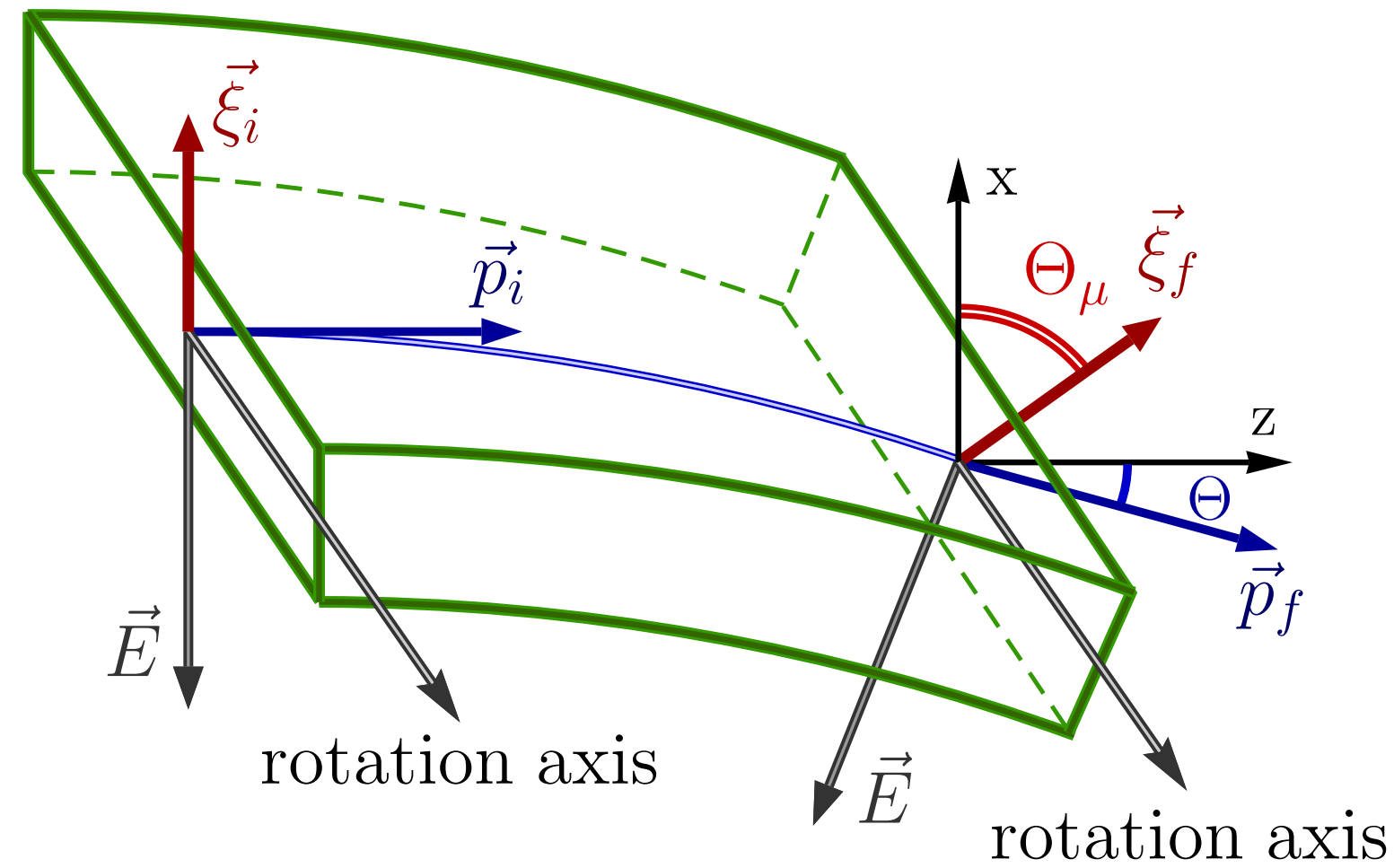
A nonzero value is forbidden by both:
T invariance and P invariance.

Particle	$ \delta $, e cm 10^{-25}
p	< 2.1
n	< 0.18
Σ^+	not measured
Λ_c^+	not measured

Spin precession in a bent crystal

V.G. Baryshevsky, Sov. Tech. Phys. Lett. 5 (1979) 73.

V.L. Lyuboshits, Sov. J. Nucl. Phys. 31 (1980) 509 [[inSPIRE](#)].



$$\Theta_\mu \equiv \angle(\xi_i \xi_f) = (1 + \gamma a) \Theta$$

$$a = \frac{g-2}{2}, \quad \Theta = \frac{L}{R}$$

γ, g, a – Lorentz factor, g -factor, anomalous MDM of Λ_c

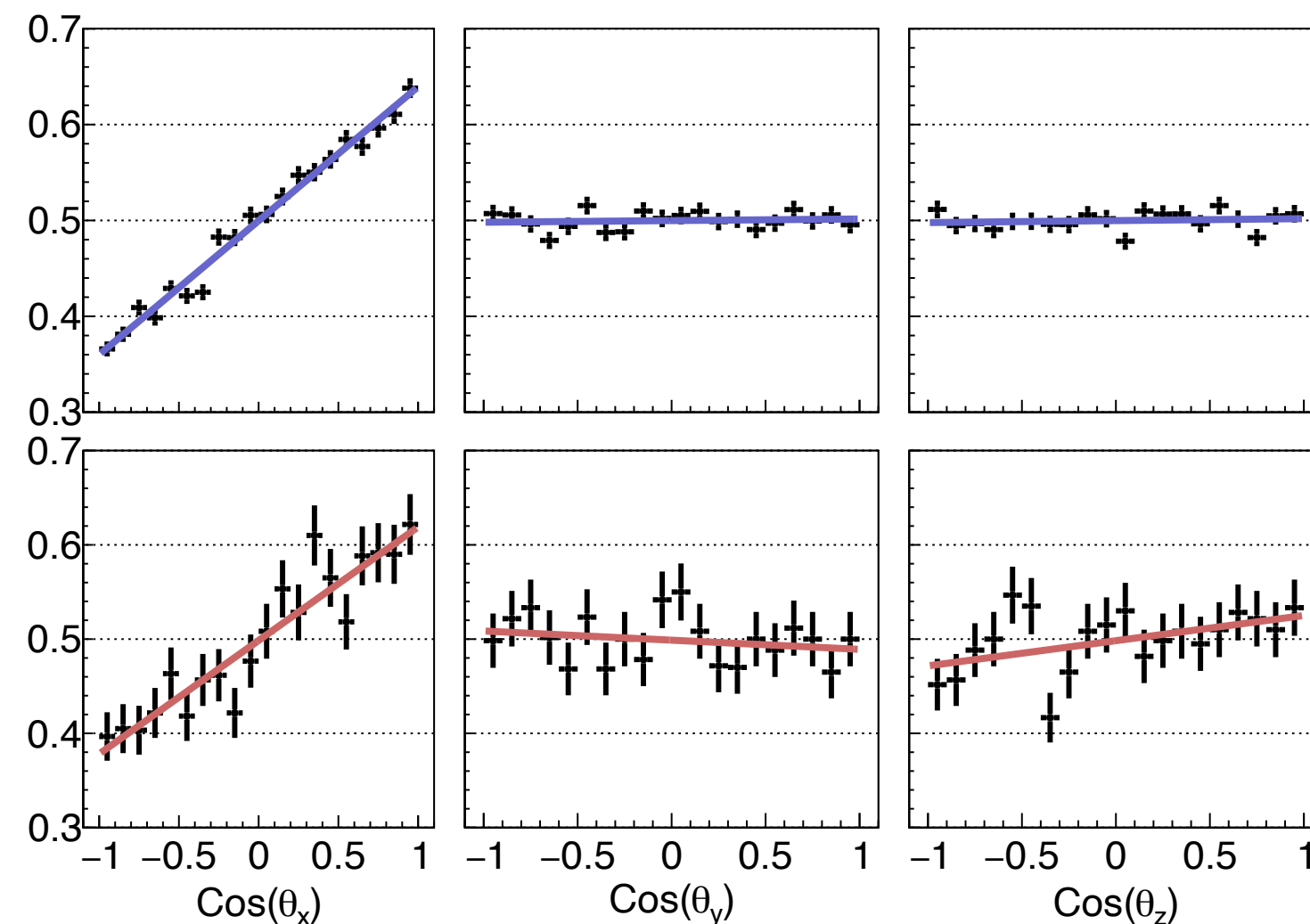
Θ, L, R – deflecting angle, length, curvature radius of the crystal

Initial Polarisation:

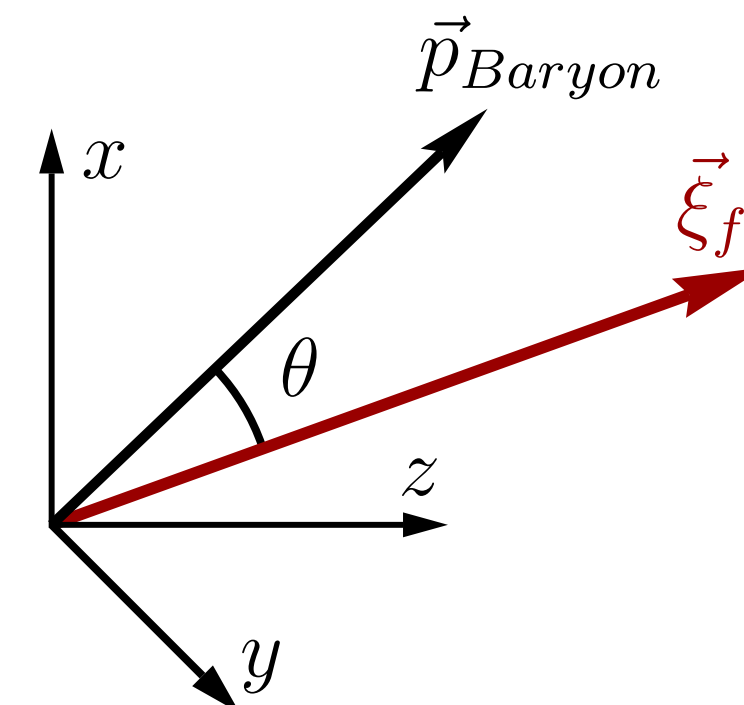
$$\vec{\xi}_i = \xi (1, 0, 0)$$

Final polarisation

$$\vec{\xi}_f = \xi (\cos \Theta_\mu, 0, \sin \Theta_\mu)$$



$\Lambda_c^+ \rightarrow \text{Meson} + \text{Baryon}$



$$\frac{dN}{d \cos \theta_z} = \frac{1}{2} \left(1 + \alpha \xi_{fz} \cos \theta_z \right)$$

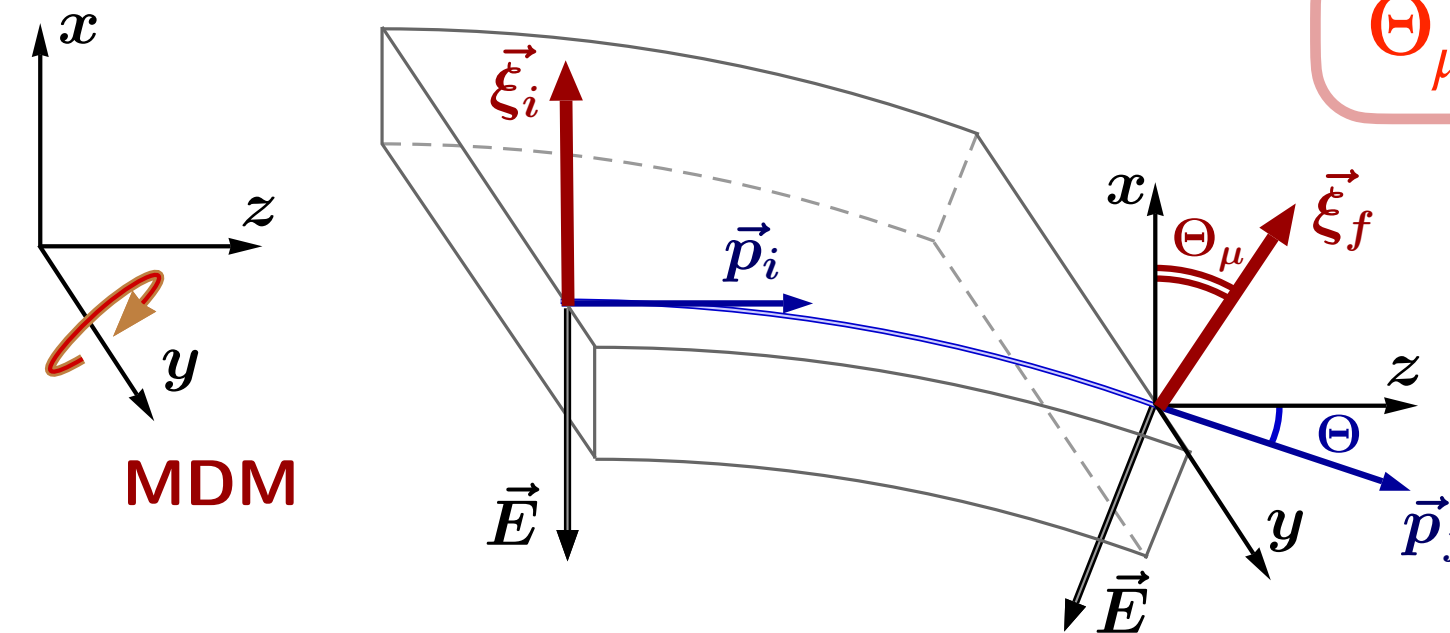
$$b \equiv \alpha \xi \Theta_\mu \quad \Delta b = \sqrt{\frac{3}{N}}$$

$$\Delta g = \frac{2}{\alpha \langle \xi \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

Optimal crystal orientation for MDM and EDM measurements

V.G. Baryshevsky,
Sov. Tech. Phys. Lett. 5 (1979) 73.

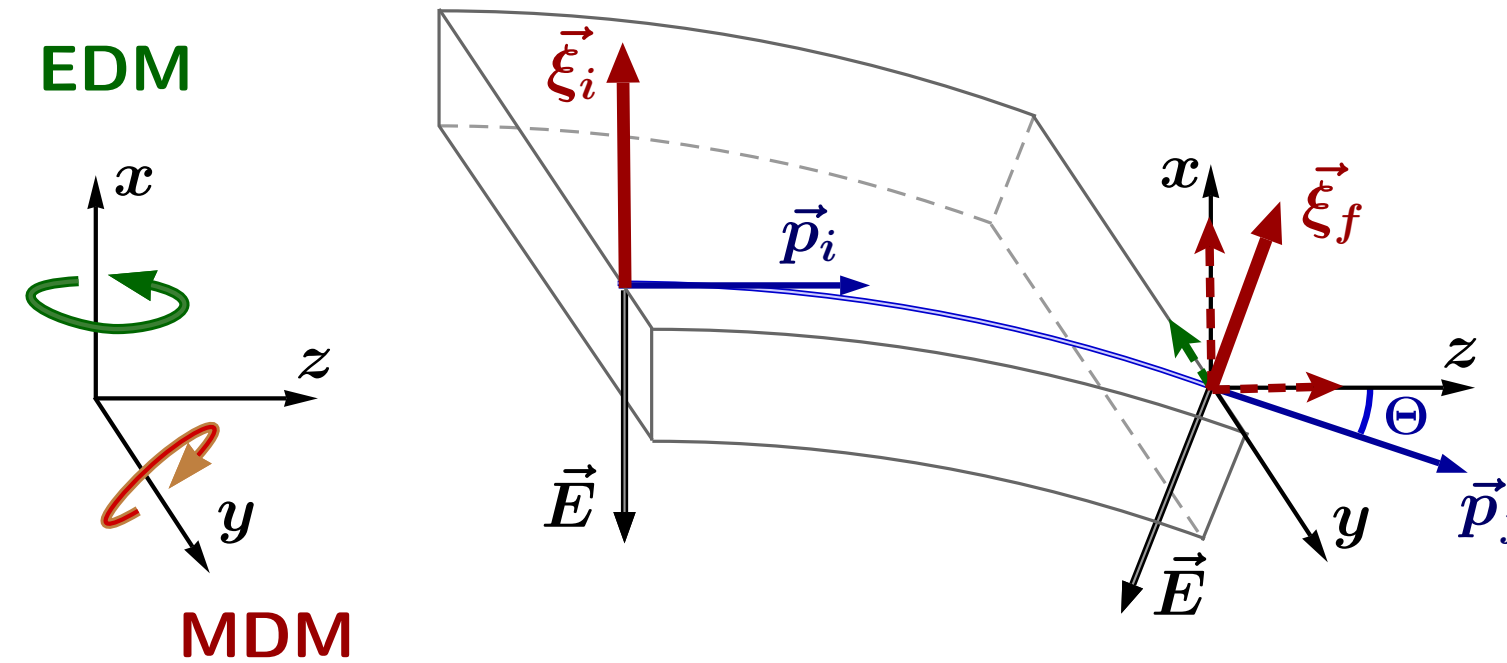
V.L. Lyuboshits,
Sov. J. Nucl. Phys. 31 (1980) 509
[\[inSPIRE\]](#).



$$\Theta_{\mu} \equiv \angle(\xi_i \xi_f) = (1 + \gamma a) \Theta$$

$$\Delta g = \frac{2}{\alpha \langle \xi_x \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

F. J. Botella et al.,
EPJ C77 (2017) 181 [\[inSPIRE\]](#)

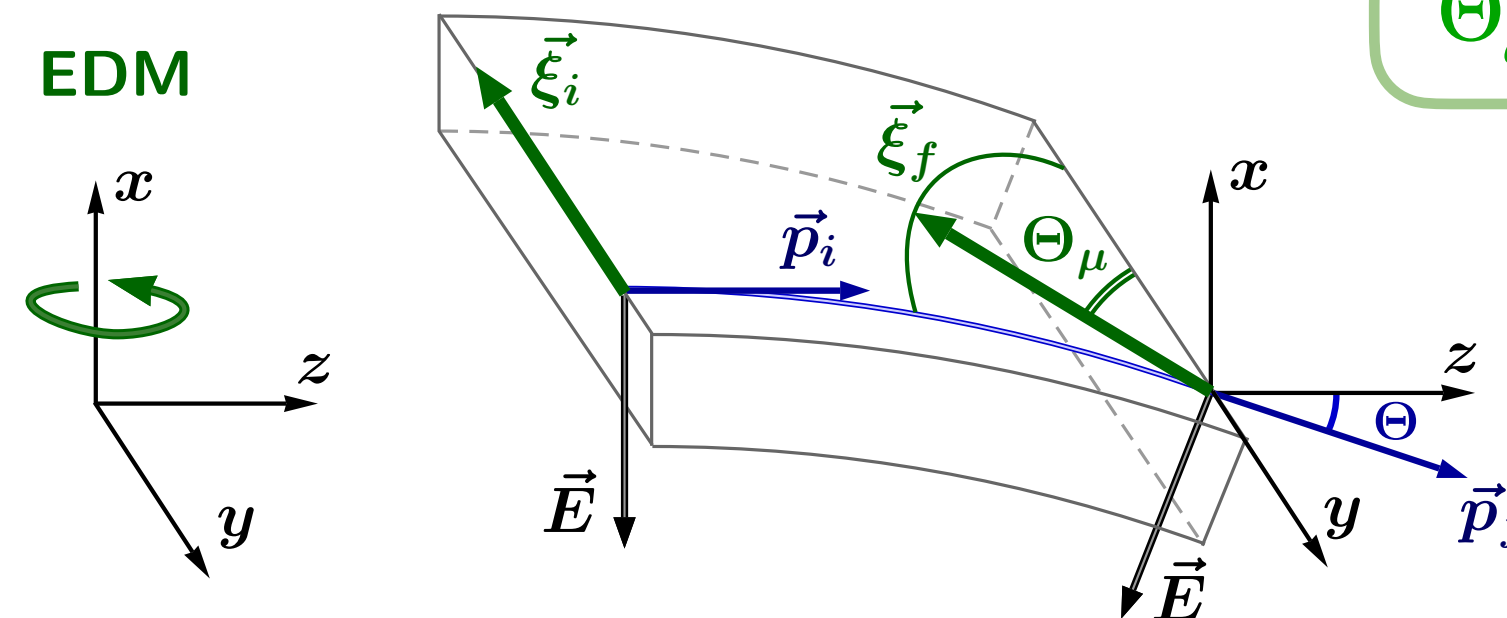


$$\frac{\Delta f}{\Delta g} = \frac{2 \gamma a}{\Theta (1 + \gamma a)^2}$$

effect is suppressed by a small bending angle

V.G. Baryshevsky,
EPJ C79 (2019) 350 [\[inSPIRE\]](#)

A.S. Fomin et al.,
EPJ C80 (2020) 358 [\[inSPIRE\]](#)



$$\Theta_d \equiv \angle(\xi_i \xi_f) = (1 + \gamma f) \Theta$$

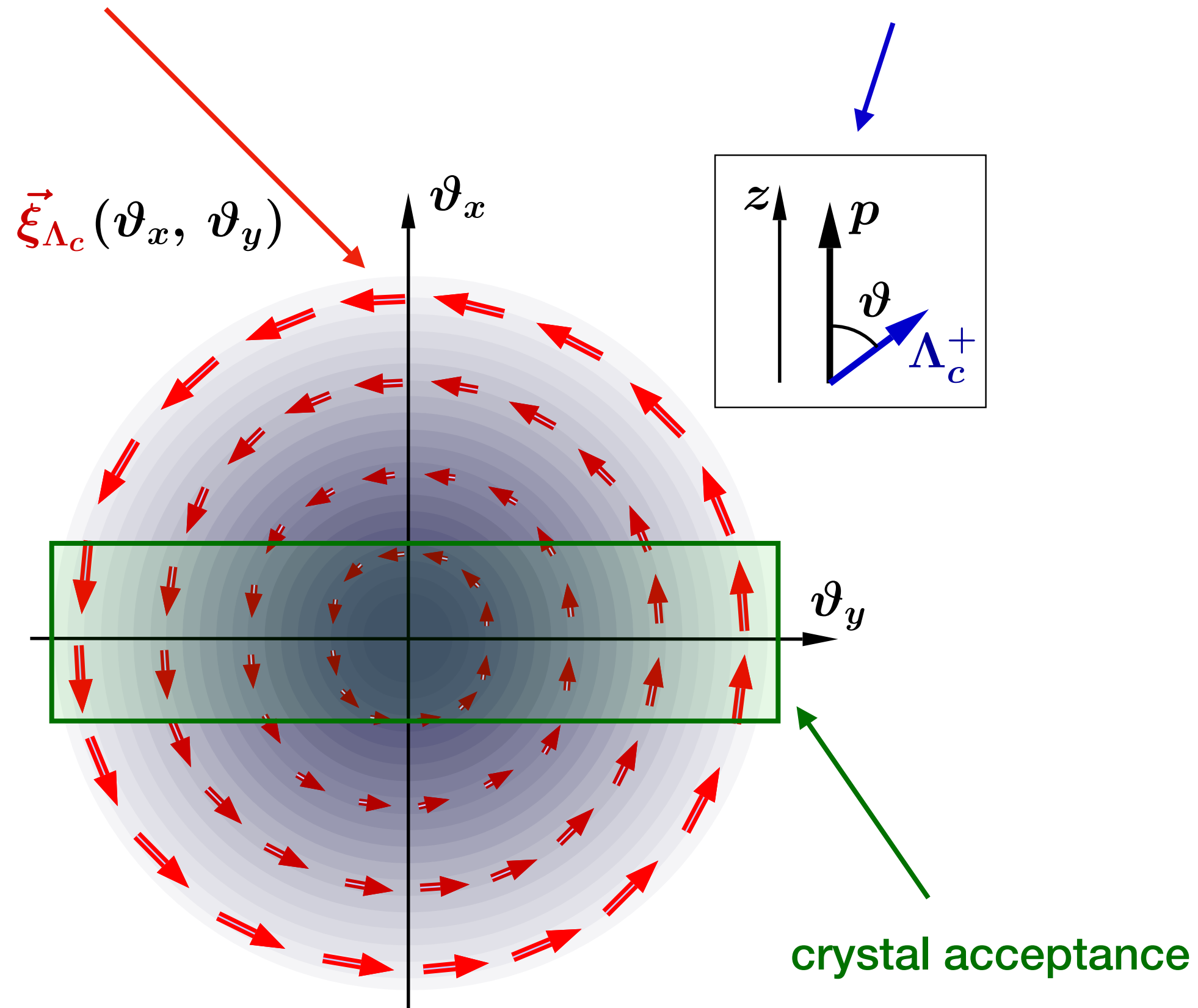
$$\Delta f = \frac{2}{\alpha \langle \xi_y \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

Optimal crystal orientation for EDM measurement: Initial polarisation

 [A. Fomin et al. Eur. Phys. J. C \(2020\) 80:358 \[1909.04654\]](#)

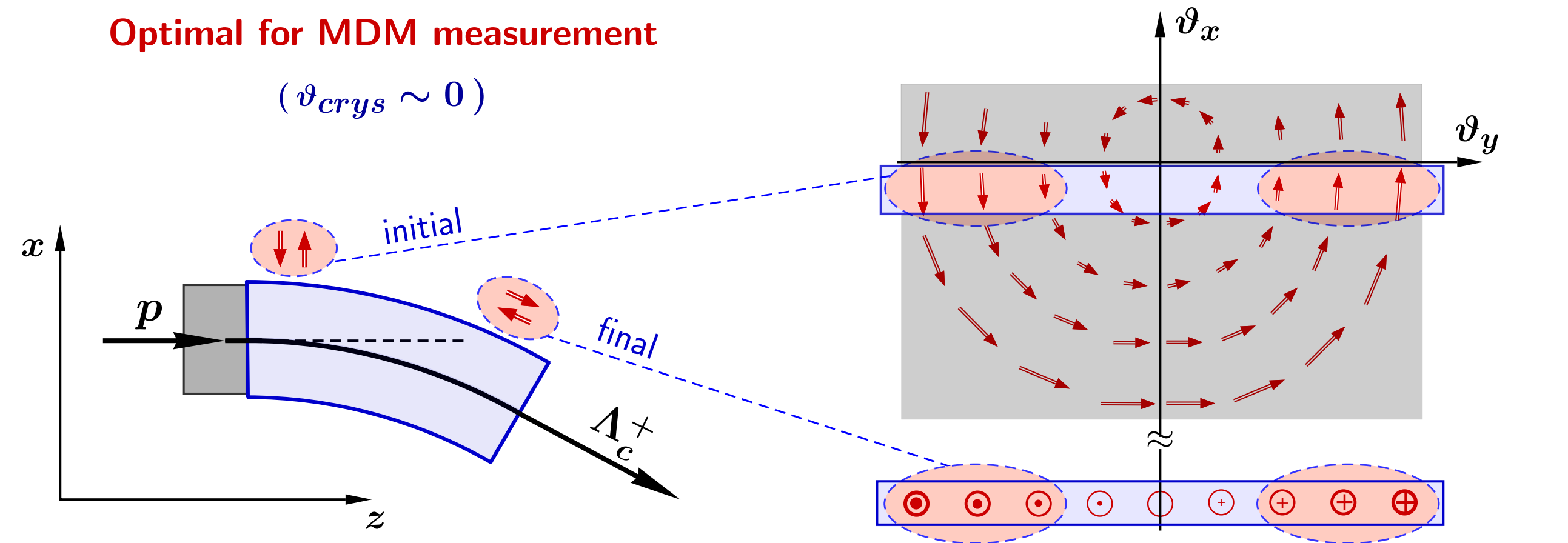
Production of Λ_c^+ in a fixed target $p + p \rightarrow \Lambda_c^+ + X$

Due to the space-inversion symmetry of the strong interaction Λ_c^+ **polarisation** is perpendicular to the **reaction plane**



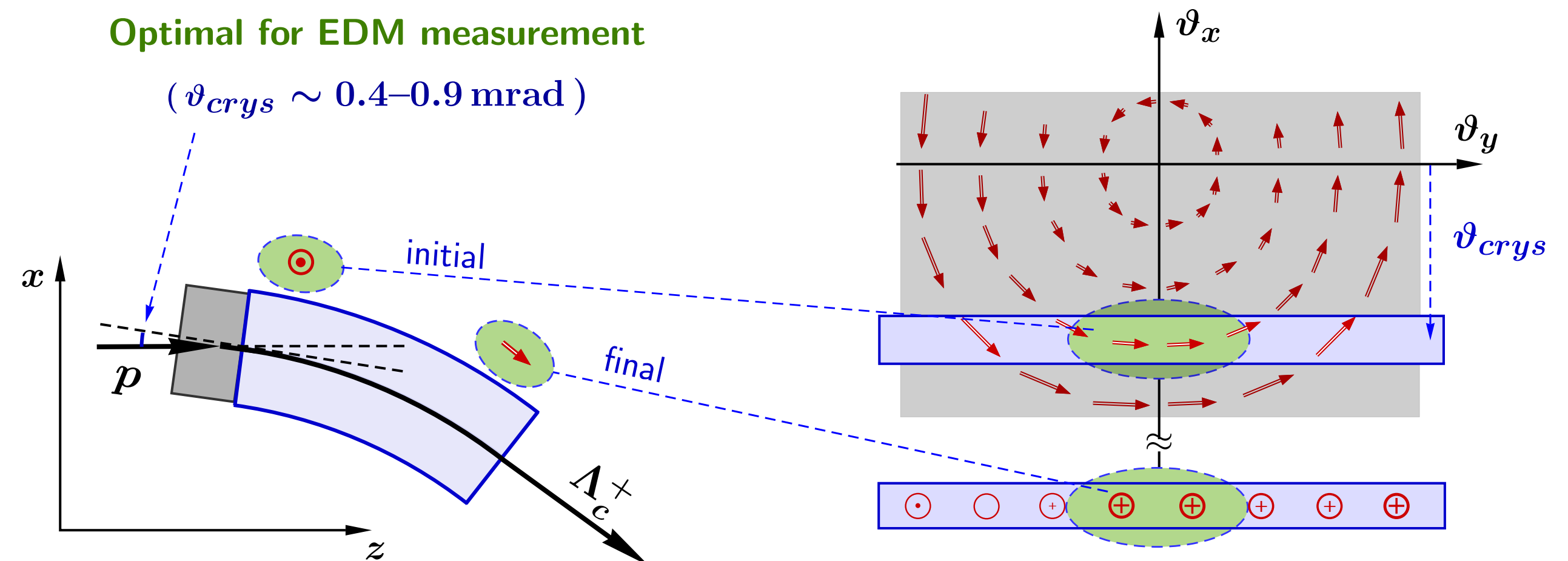
Optimal for MDM measurement

$(\vartheta_{crys} \sim 0)$



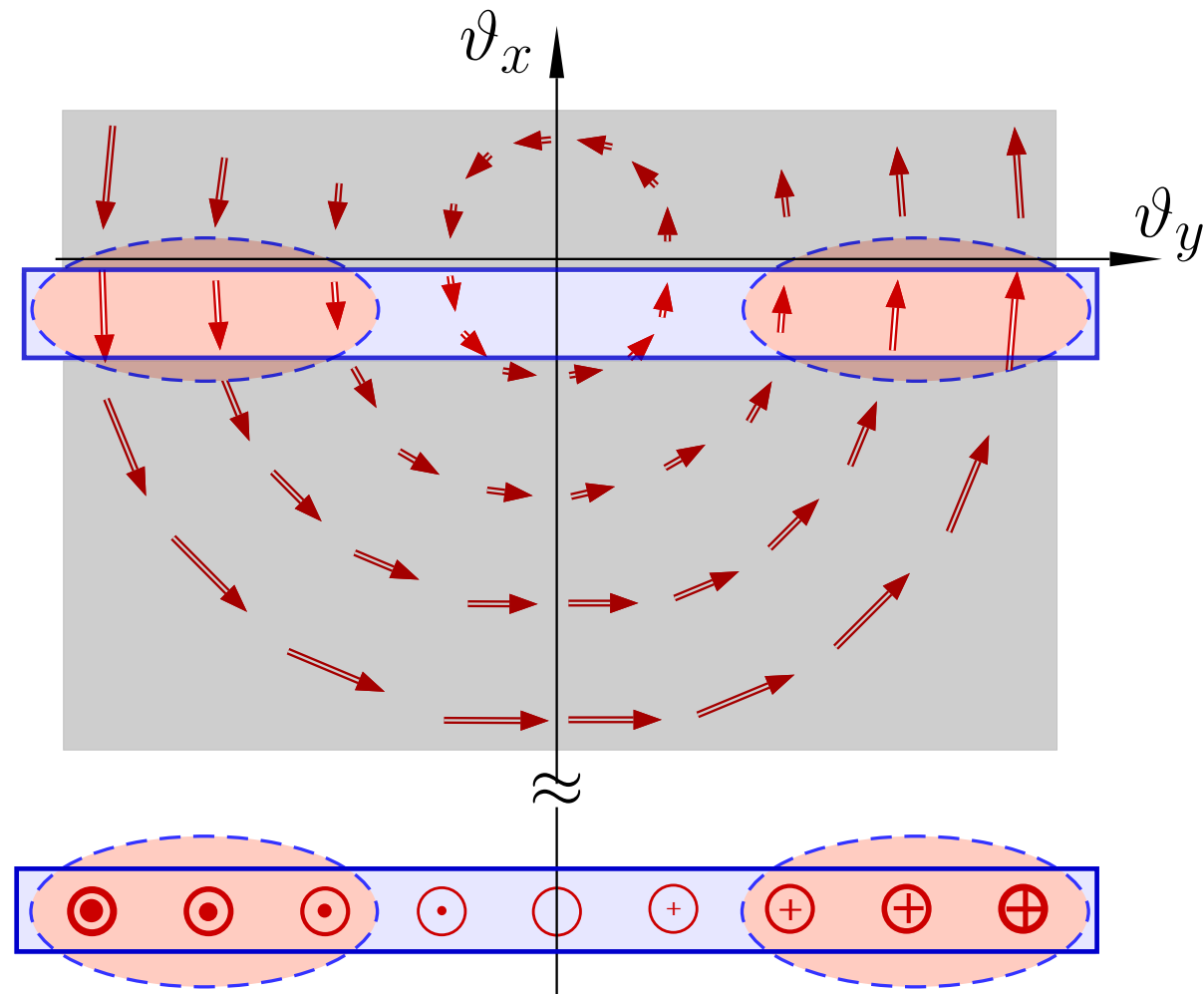
Optimal for EDM measurement

$(\vartheta_{crys} \sim 0.4-0.9 \text{ mrad})$

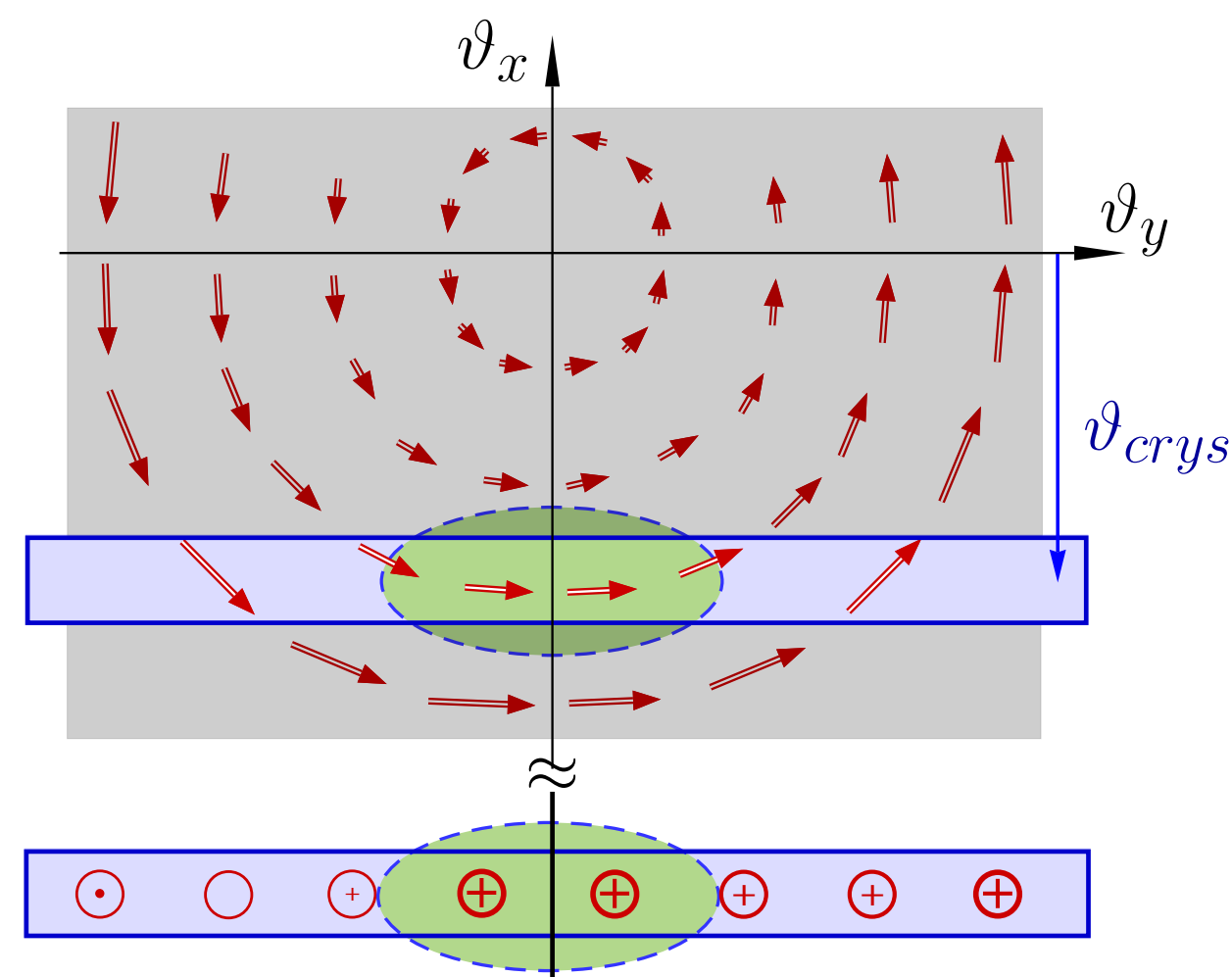


Optimal crystal orientation for EDM measurement: Quantitative analysis

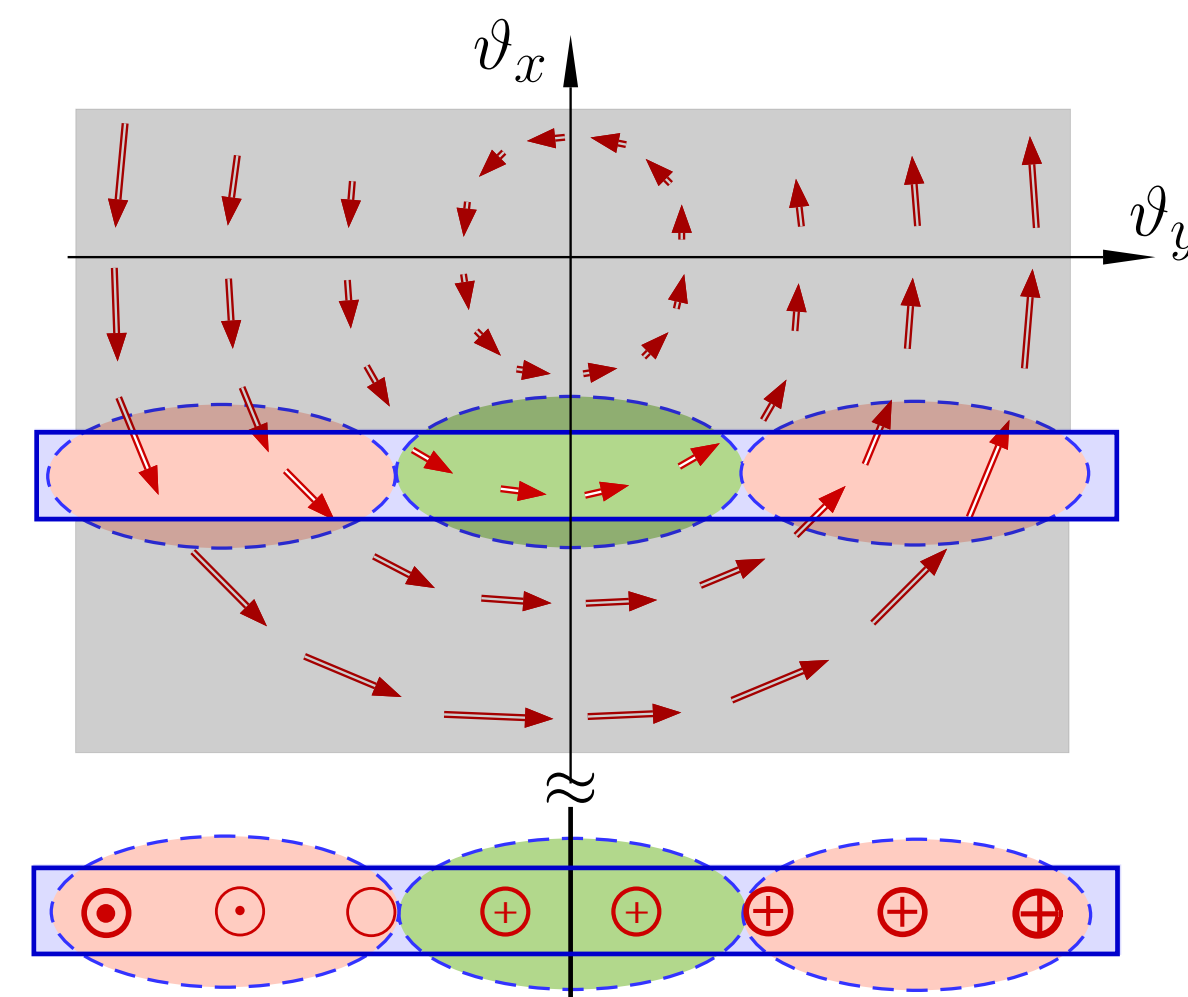
Optimal for MDM measurement



Optimal for EDM measurement



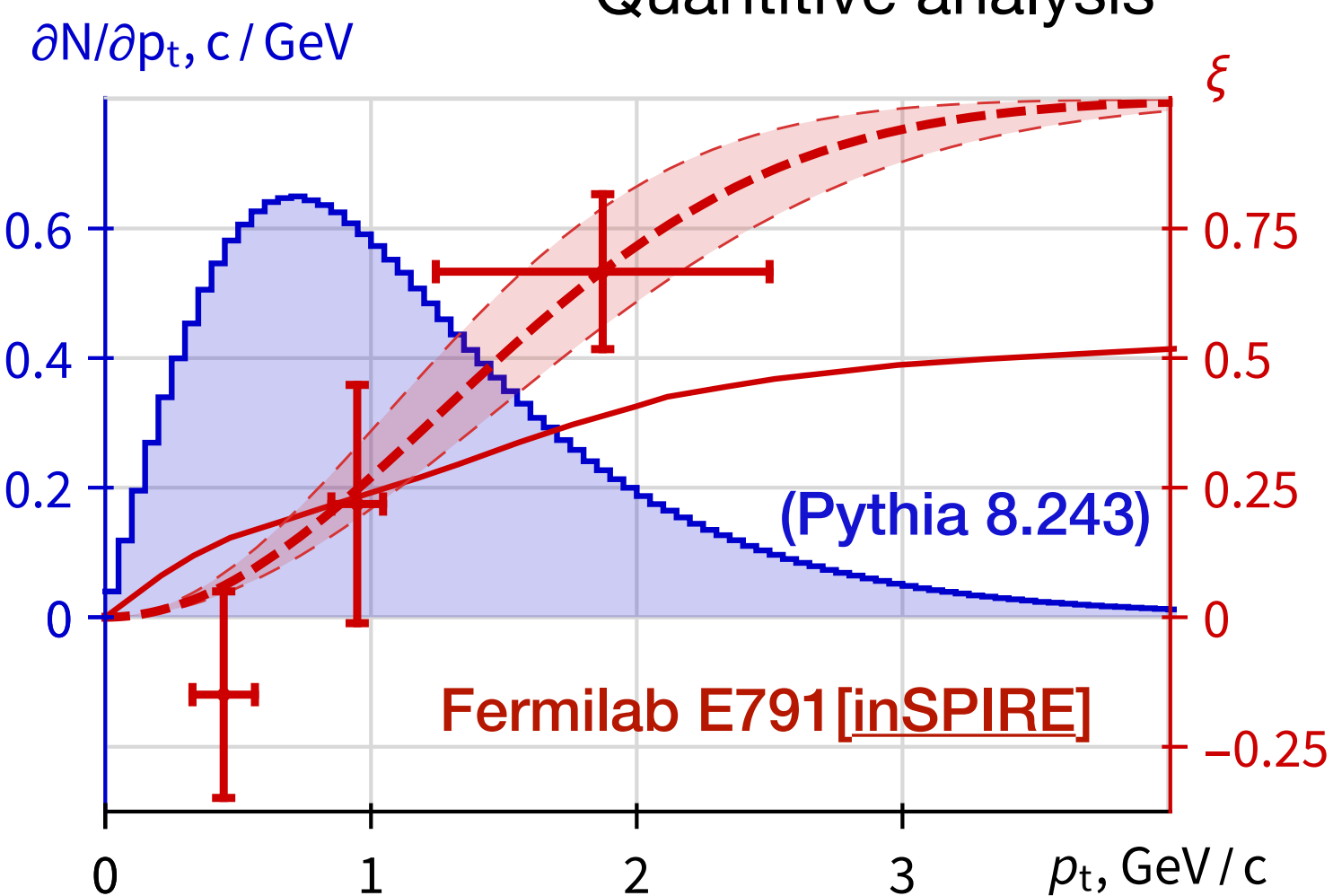
Simultaneous measurement



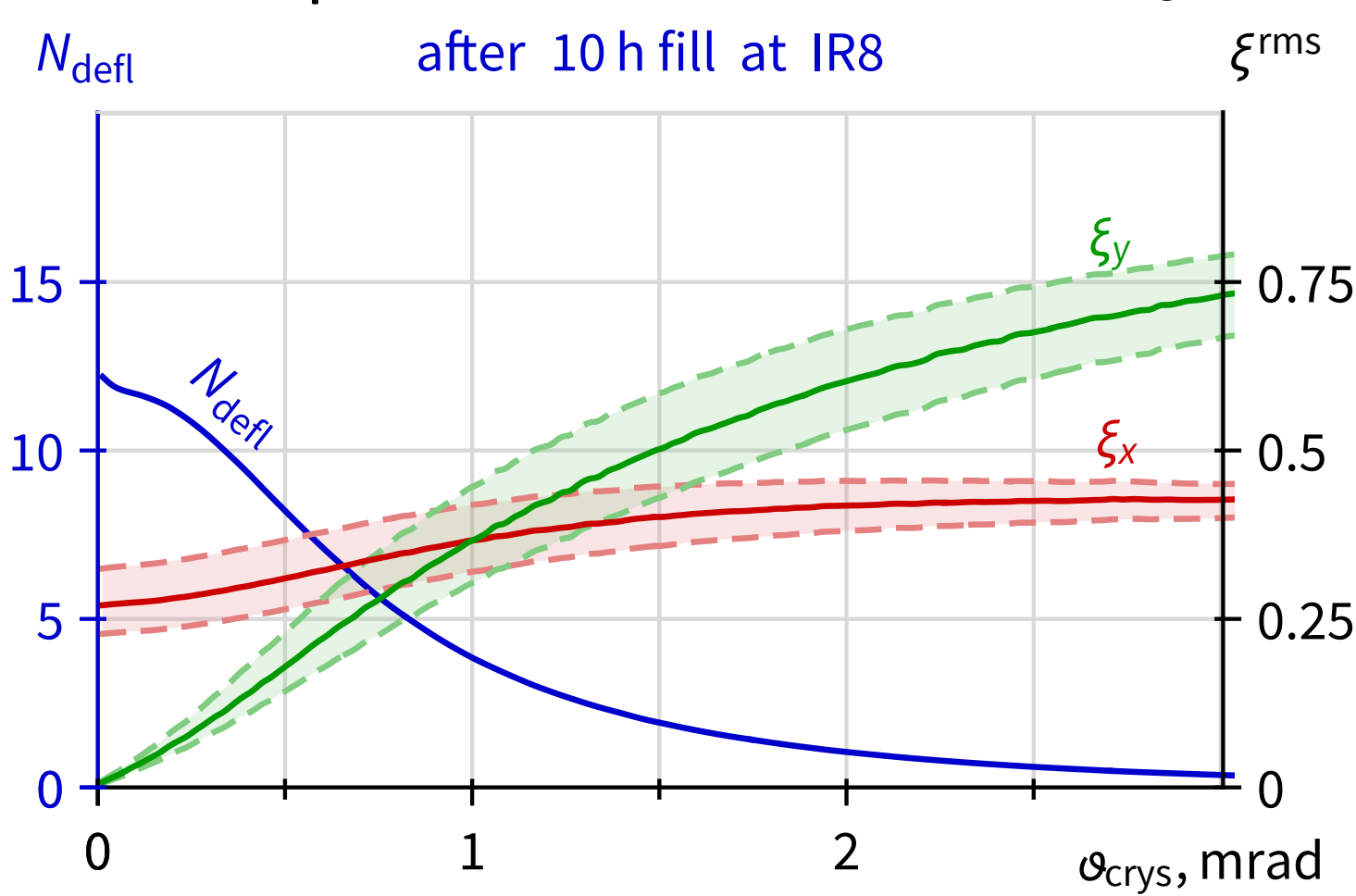
$$\Delta g = \frac{2}{\alpha \langle \xi_x \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

$$\Delta f = \frac{2}{\alpha \langle \xi_y \gamma \rangle \Theta} \sqrt{\frac{3}{N}}$$

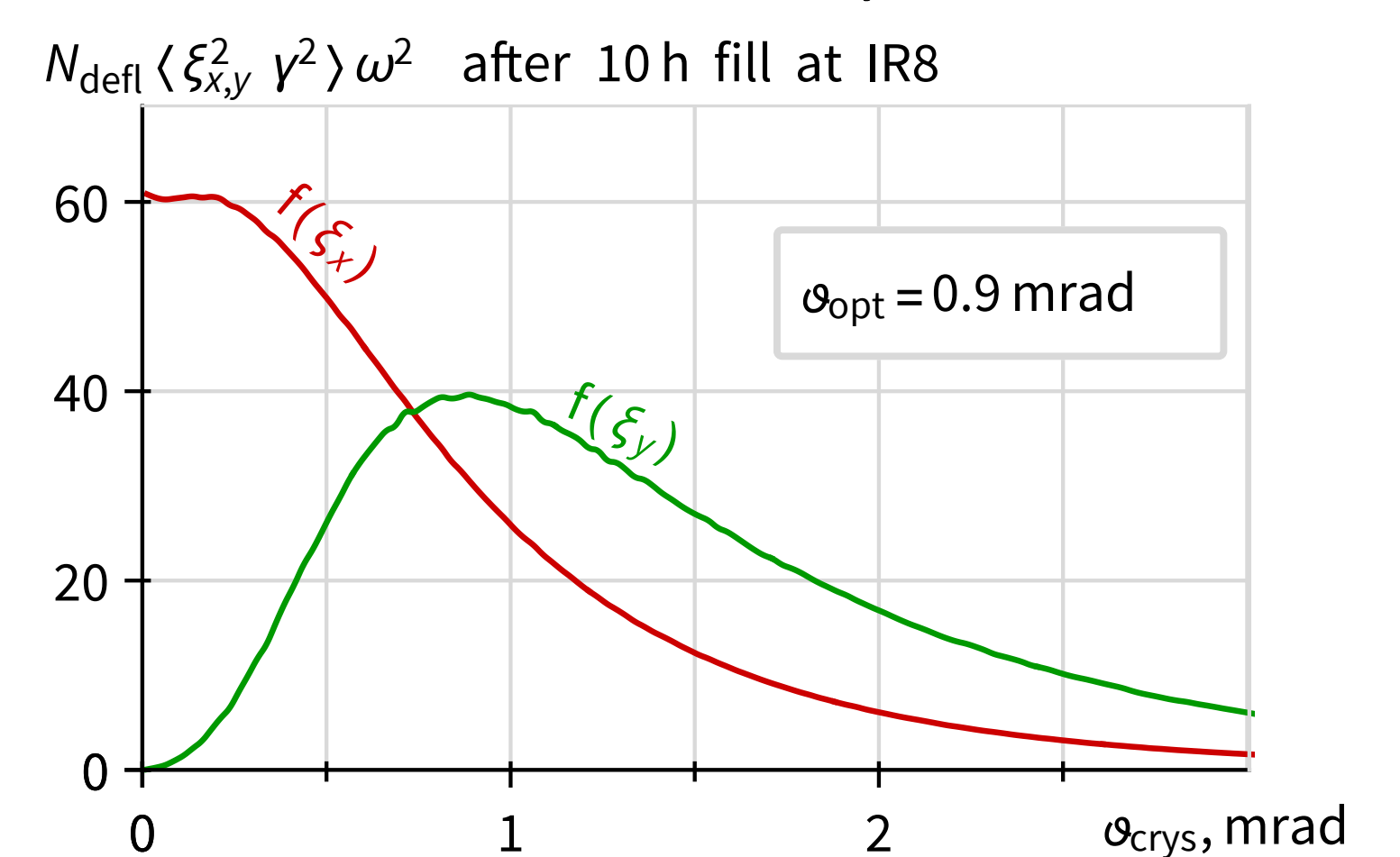
Quantitative analysis



Initial polarisation of deflected Λ_c^+

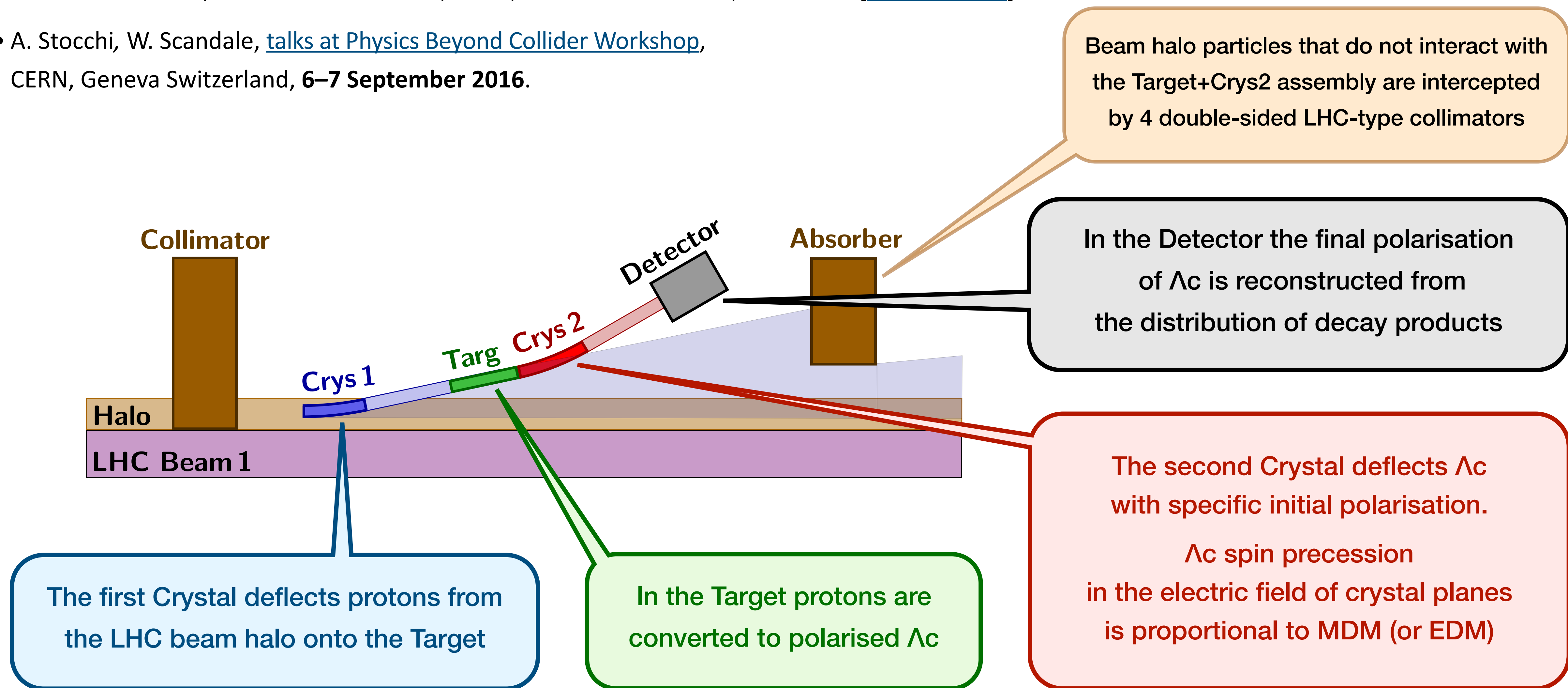


Measurement efficiency



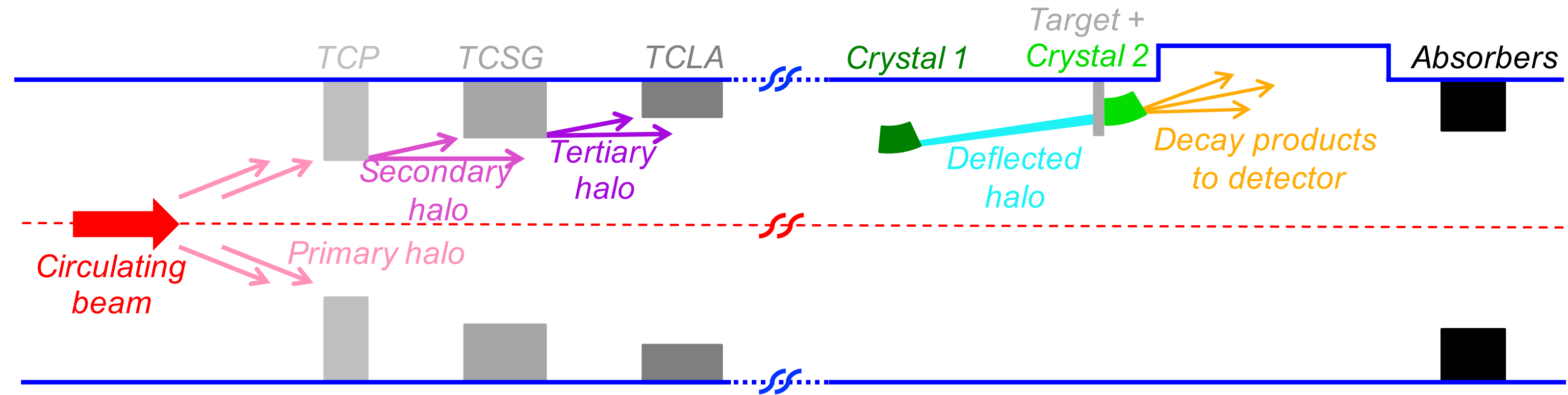
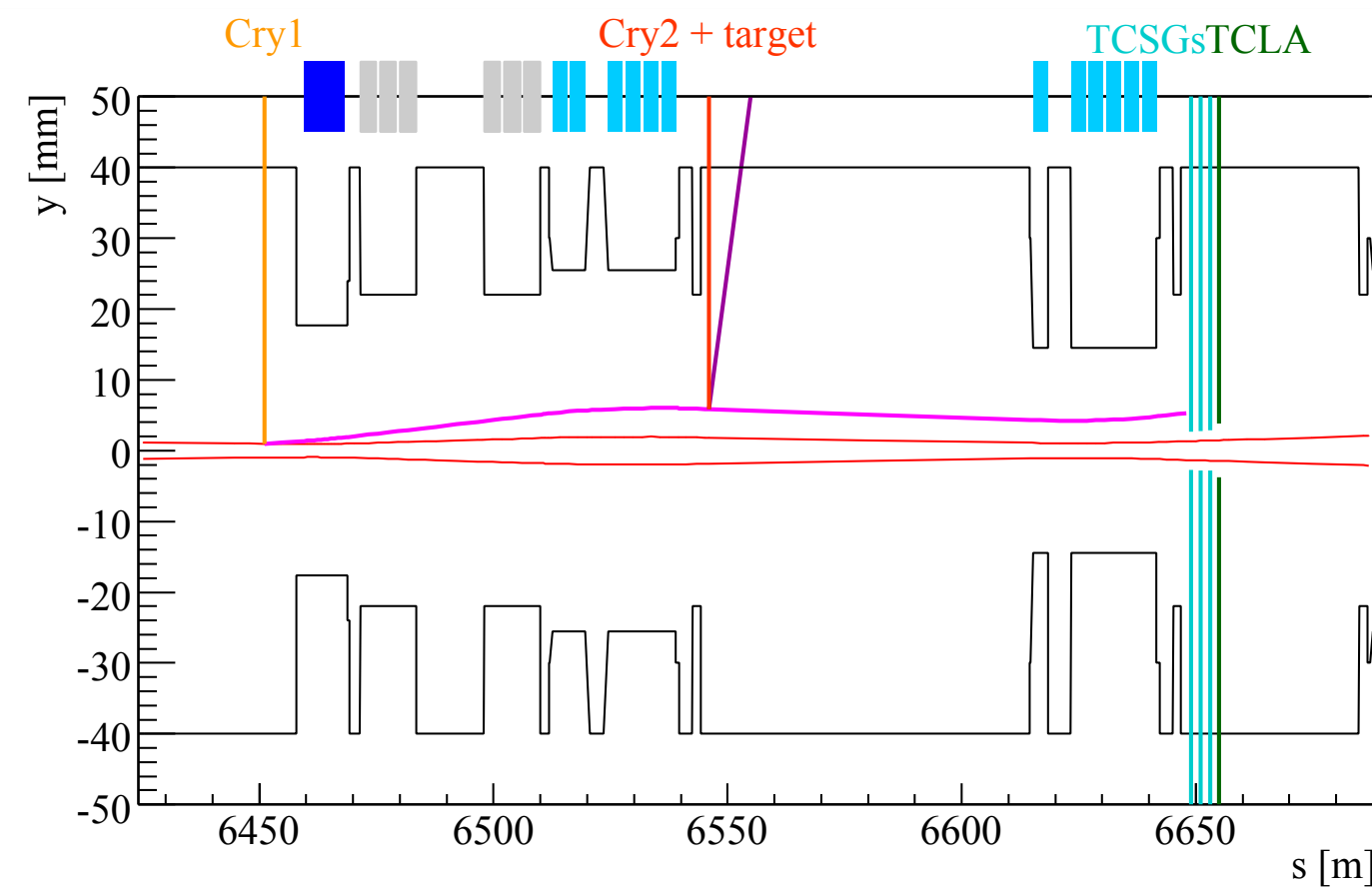
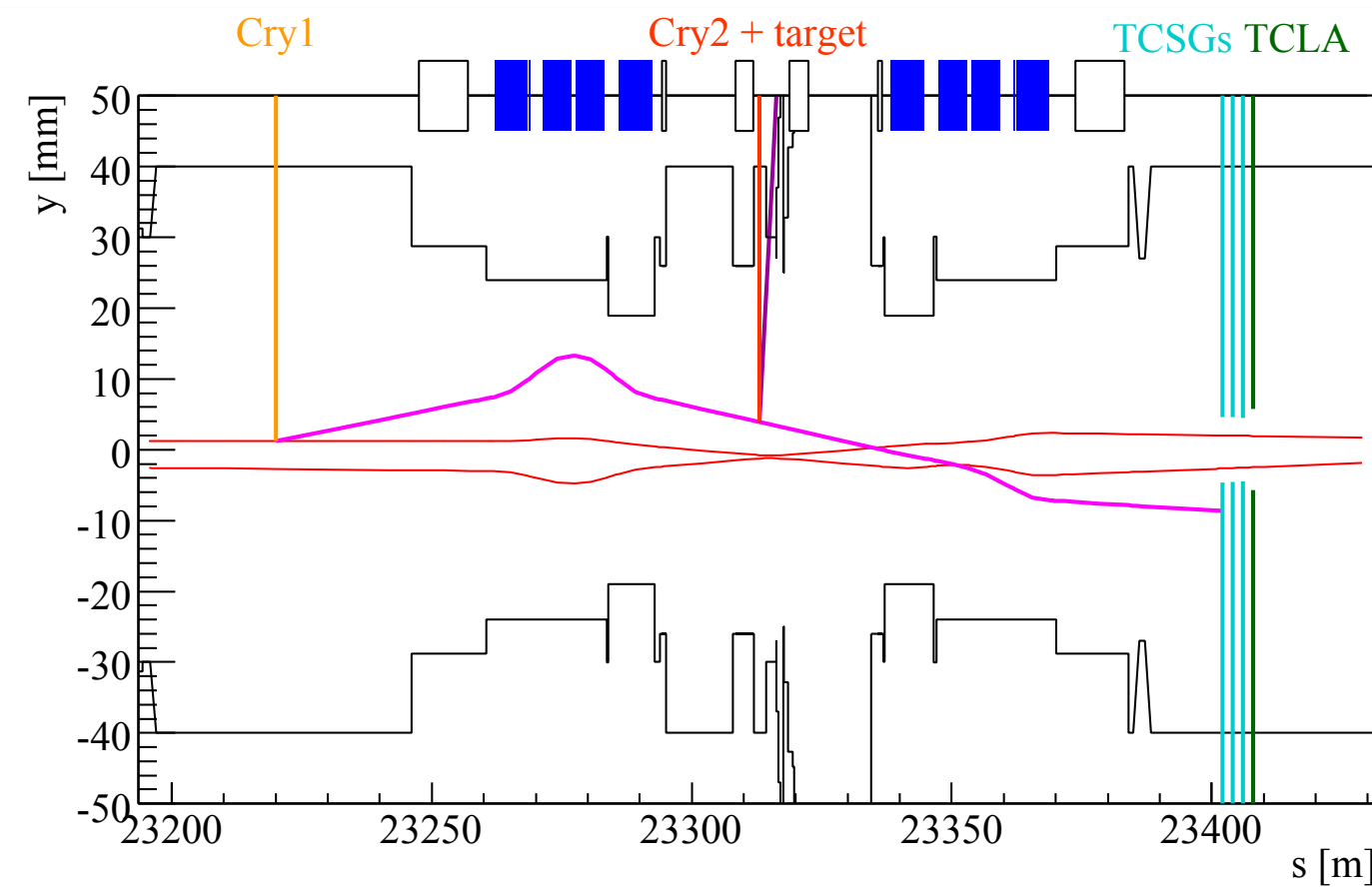
MDM and EDM of charmed baryons: Fixed target at the LHC

- L. Burmistrov et al., CERN-SPSC-2016-030, CERN, Geneva Switzerland, **June 2016** [[SPSC-EOI-012](#)].
- A. Stocchi, W. Scandale, [talks at Physics Beyond Collider Workshop](#), CERN, Geneva Switzerland, **6–7 September 2016**.

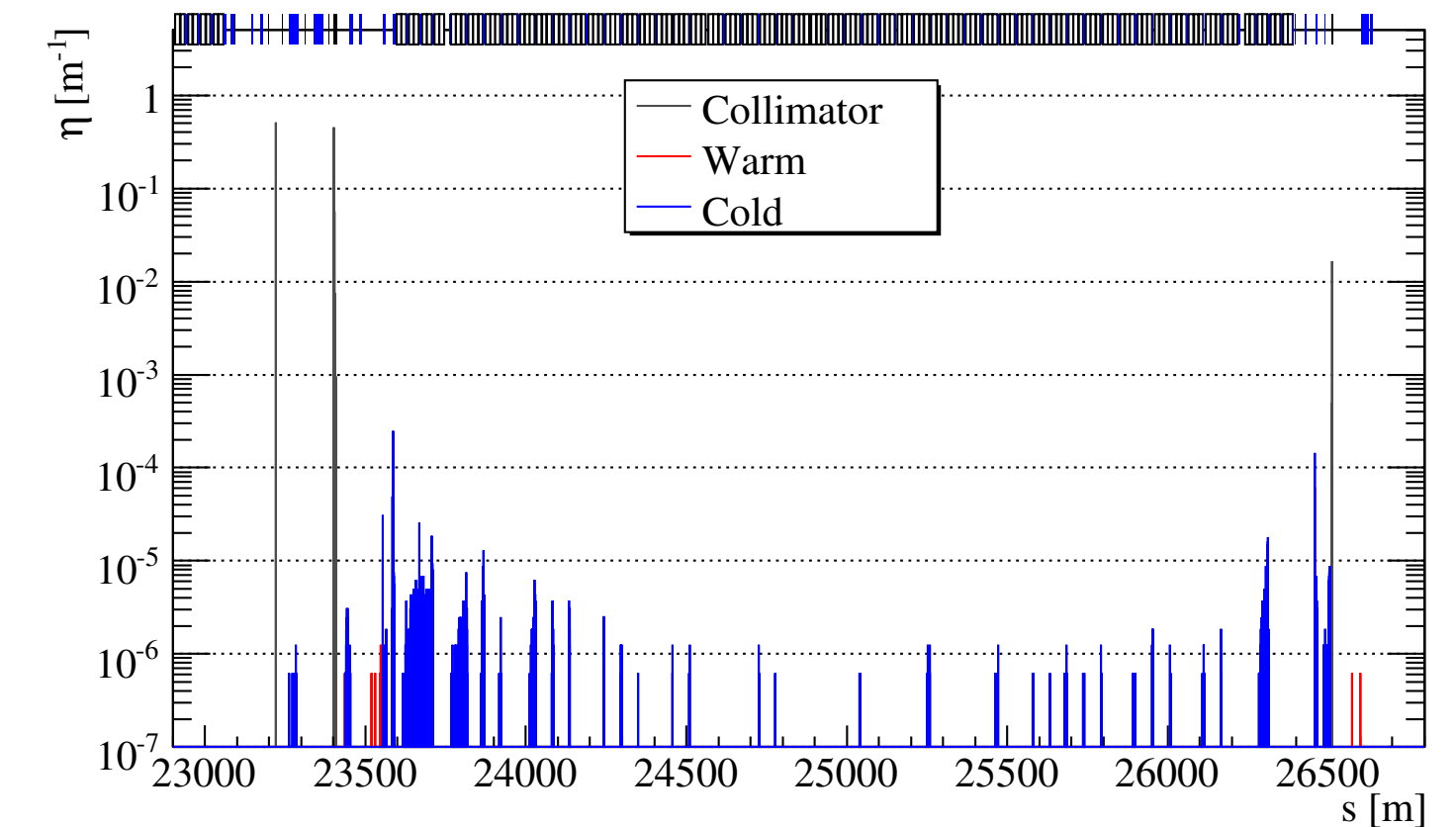


Performance assessment of layouts in IR3 and IR8 of LHC

D. Mirarchi et al. Eur. Phys. J. C 80 (2020) 10, 929



- impact on the machine
- optimisation of Crystal 1 and Absorbers positions
- running experiment in a parasitic mode
- layout in front of LHCb (IR8) 4.3×10^{10} POT/fill
- alternative layout at IR3 3.0×10^{10} POT/fill
- restriction on Crystal 2 bending radius

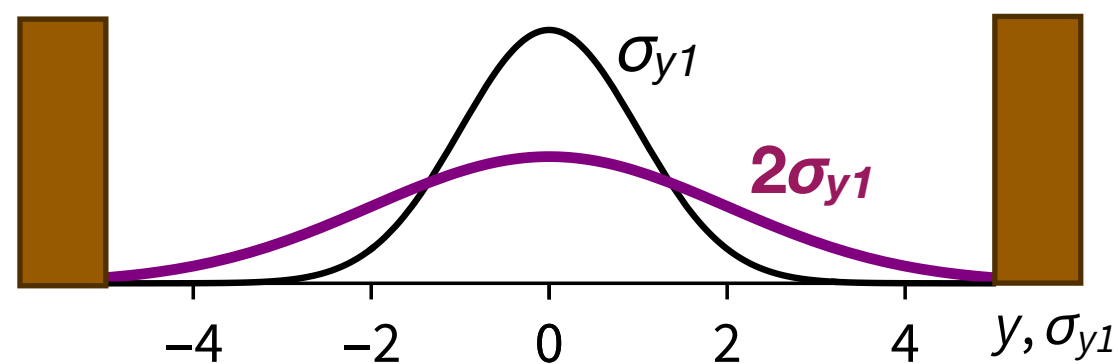


Increasing the PoT through active bunch excitation

General idea

- apply random noise excitations with the **transverse damper** when a selected trains of bunches pass by

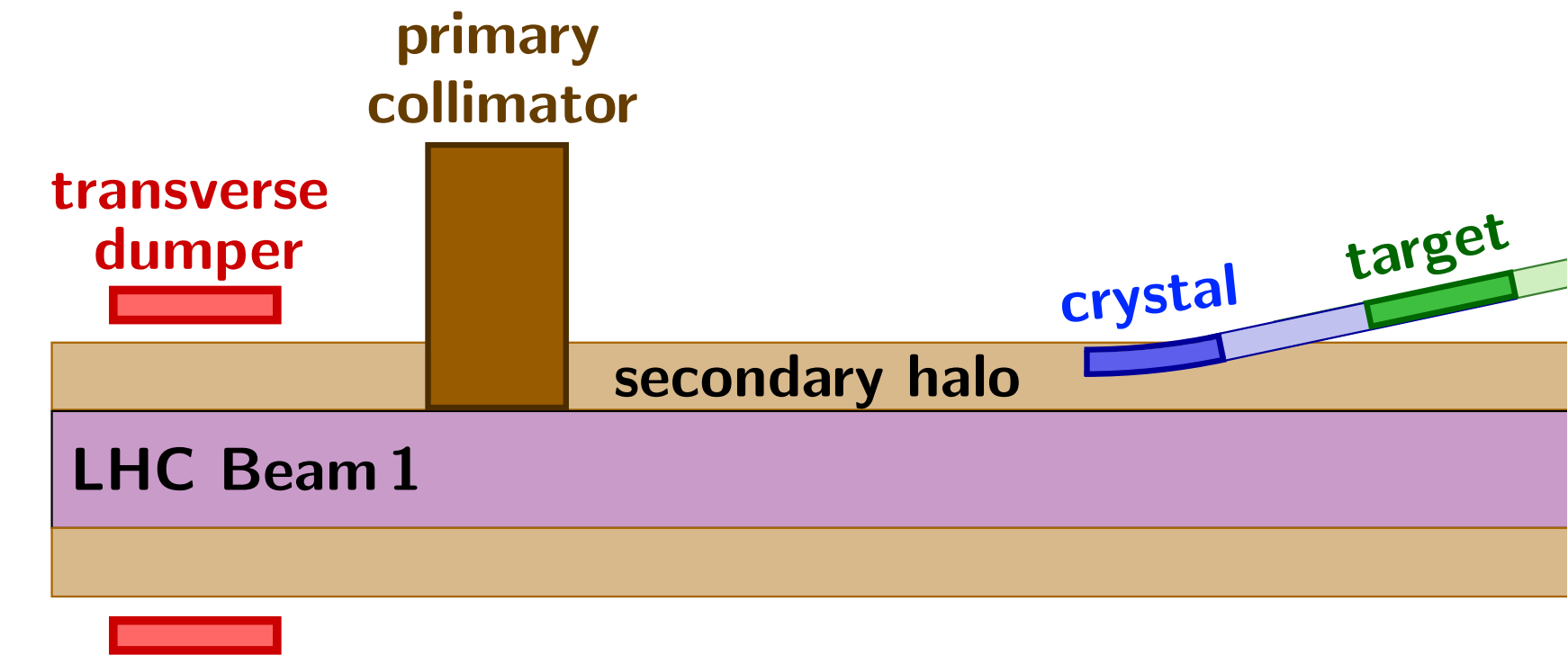
→ **emittance grow** of these bunches



→ increase the loss rate in **primary collimators** at IR7

→ enrich the **secondary halo** of the beam

→ increase the flux on the **crystal** → increase the deliverable rates of **PoT**

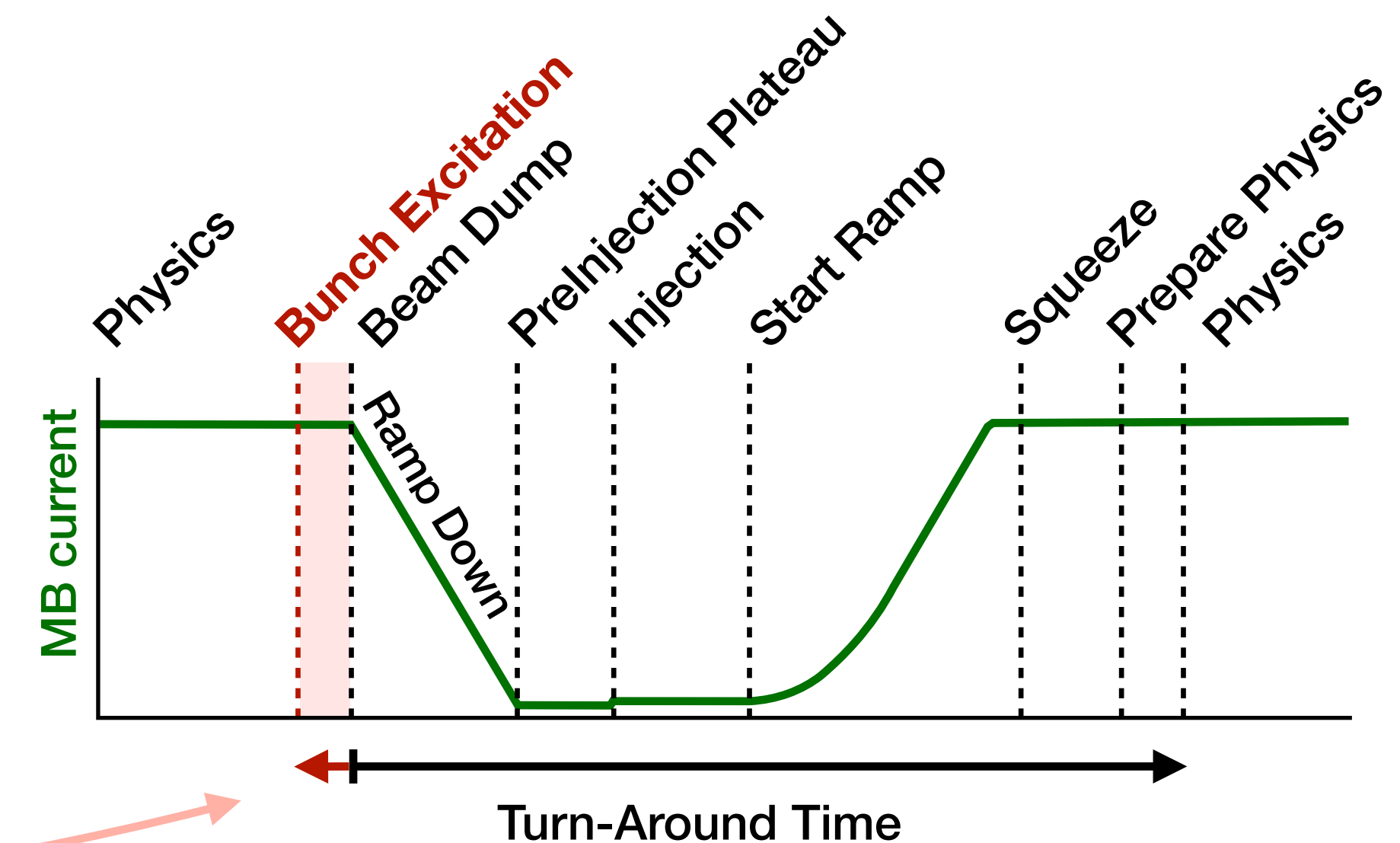
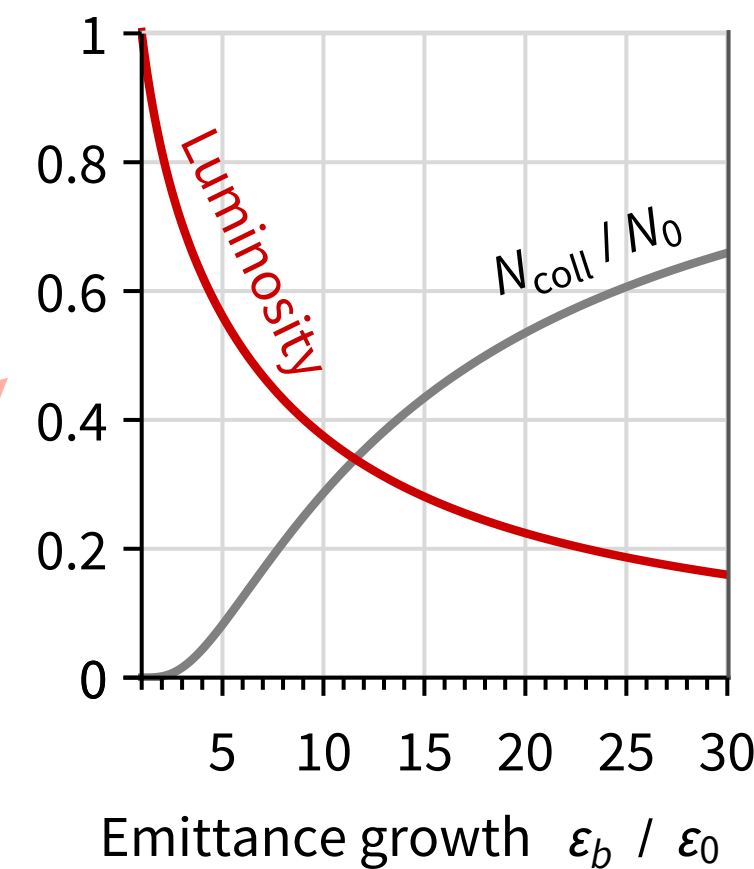


Transverse damper is currently used in LHC

- corrects the trajectory of the bunches during the fill
- increase the loss rate during alignment of collimators

Drawback — reduction of total luminosity

- excitation during the fill → higher emittance
- excitation before dumping → longer turn-around time



Maximum achievable PoT due to beam excitation (based on 2018 run data)

The **gain factor** in PoT from the beam excitation with respect to the “Baseline” (no excitation):

$$\eta_{\text{exc}} = \frac{N_{\text{PoT}}^{(\text{ex})}}{N_{\text{PoT}}^{(\text{BL})}}$$

too high rates of PoT per bunch crossing (~49 @LHCb)

200 kW in coll. system		3.5 p / bunch X @ LHCb		
after collisions			after collis.	during collis.
t, h	$\Delta L / L$	t, h	$\Delta L / L$	$\Delta L / L$
1908	0	1908	0	0
8	$> 0.4 \%$	167	8.7 %	2.9 %
-	-	54	2.8 %	1.2 %

The most efficient scenario:

The gradual excitation of bunches (3.5 p/bunch X) at the end of the fills with 2556 bunches, with duration under 1 hour:

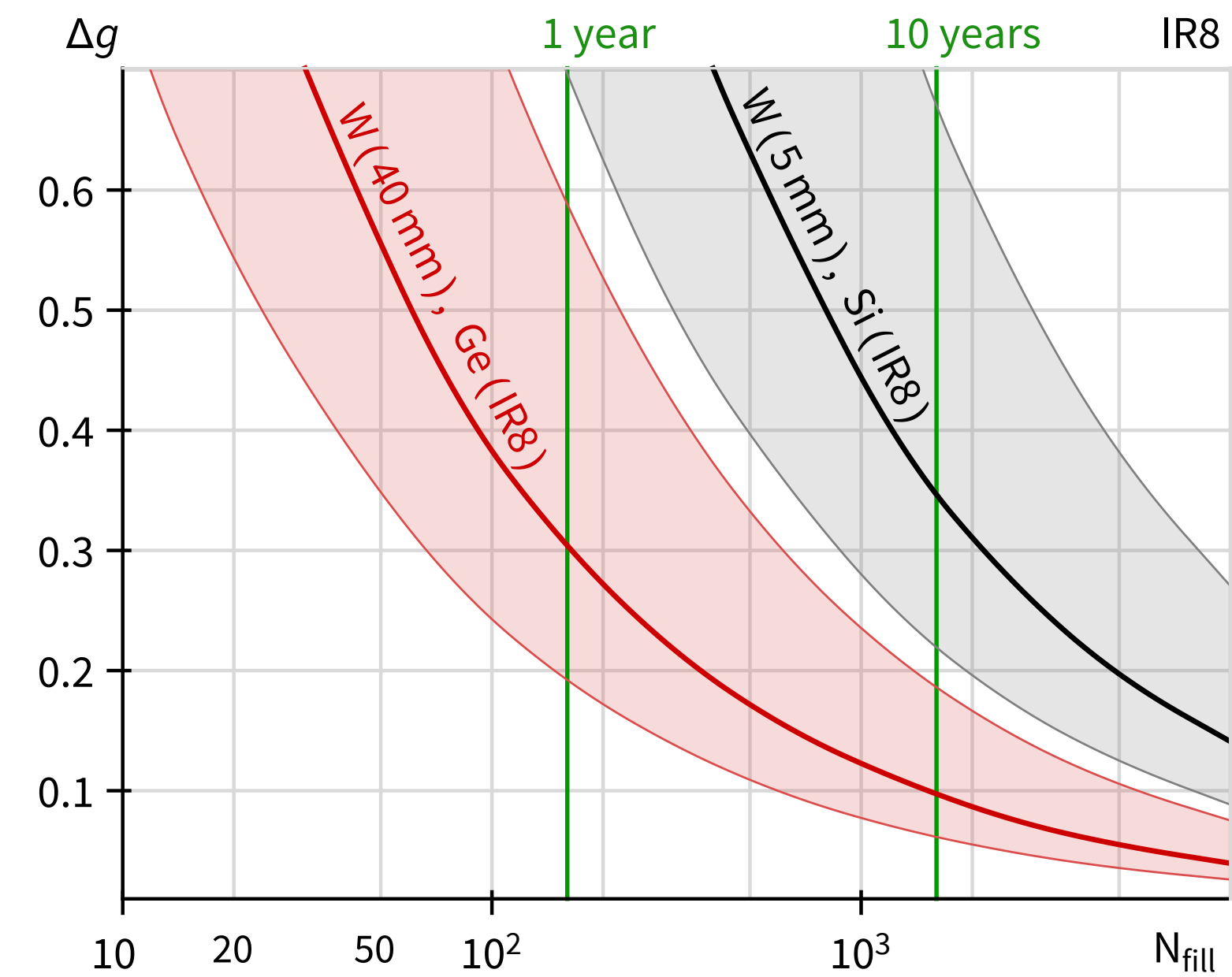
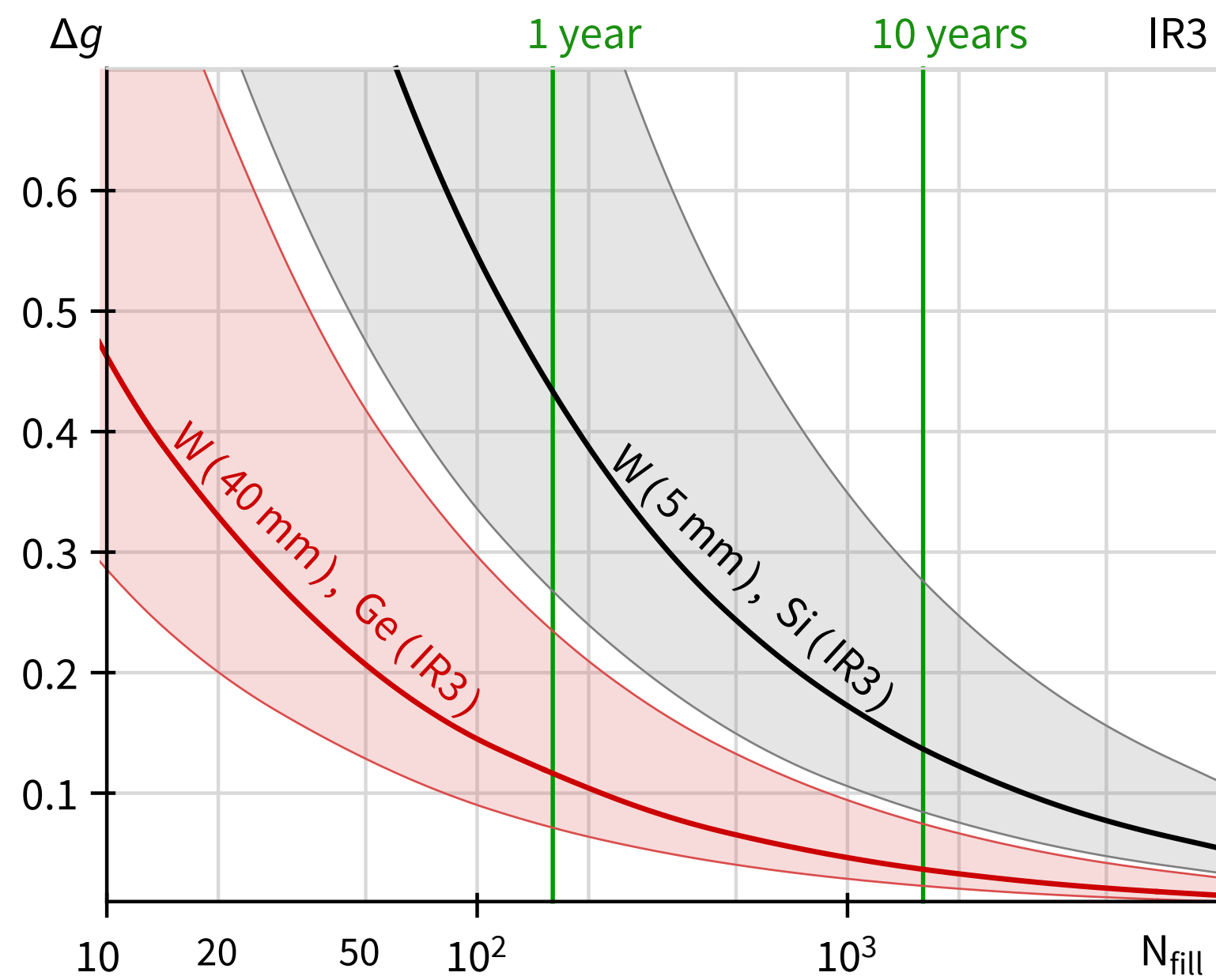
⇒ delivers **3 times more PoT** w.r.t. “Baseline”

⇒ reduces the total luminosity at ATLAS and CMS by **~1.2%**

Performance assessment of layouts in IR3 and IR8: possible improvements

 [A. Fomin et al. EPJ C80 \(2020\) 358](#)

- Thicker target 5 mm → 40 mm:
ionisation energy losses and
multiple scattering can be neglected,
showers production - to be checked
- Proton rate, $3\text{--}4.3\times10^{10}$ per 10h fill
[D. Mirarchi et al. EPJ C80 \(2020\) 10, 929](#)



Possible improvements:

	1 → 2	t1 / t2
Target	5 mm → 40 mm	4
Crystal	silicon → germanium	4
Detector	LHCb (IR8) → dedicated at IR3	5
Beam excitation	additional to parasitic scenario	3

- **10 years** at LHCb, $\sim 7\times10^{13}$ POT, 5mm, Si → $\Delta g \sim \mathbf{0.35}$
- **1 year** at IR3, $\sim 0.5\times10^{13}$ POT, 40mm, Ge → $\Delta g \sim \mathbf{0.12}$
- **10 years** at LHCb, $\sim 7\times10^{13}$ POT, 40mm, Ge → $\Delta d \sim \mathbf{2.6\ 10^{-16}\ e\ cm}$
(optimal orientation → data taking time reduced by **~170**)
- big uncertainty ($\times 10$) due to α parameter
→ significant improvement in PhD thesis of Elisabeth Neil

thank you